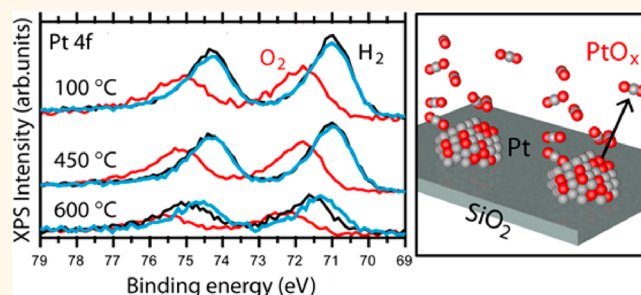


# Stability of Platinum Nanoparticles Supported on SiO<sub>2</sub>/Si(111): A High-Pressure X-ray Photoelectron Spectroscopy Study

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**ABSTRACT** The stability of Pt nanoparticles (NPs) supported on ultrathin SiO<sub>2</sub> films on Si(111) was investigated *in situ* under H<sub>2</sub> and O<sub>2</sub> (0.5 Torr) by high-pressure X-ray photoelectron spectroscopy (HP-XPS) and *ex situ* by atomic force microscopy (AFM). No indication of sintering was observed up to 600 °C in both reducing and oxidizing environments for size-selected Pt NPs synthesized by inverse micelle encapsulation. However, HP-XPS revealed a competing effect of volatile PtO<sub>x</sub> desorption from the Pt NPs (~2 and ~4 nm NP sizes) at temperatures above 450 °C in the presence of 0.5 Torr of O<sub>2</sub>. Under oxidizing conditions, the entire NPs were oxidized, although with no indication of a PtO<sub>2</sub> phase, with XPS binding energies better matching PtO. The stability of catalytic NPs in hydrogenation and oxidation reactions is of great importance due to the strong structure sensitivity observed in a number of catalytic processes of industrial relevance. An optimum must be found between the maximization of the surface active sites and metal loading (*i.e.*, minimization of the NP size), combined with the maximization of their stability, which, as it will be shown here, is strongly dependent on the reaction environment.



**KEYWORDS:** platinum nanoparticle · inverse micelle · volatile PtO<sub>x</sub> · high-pressure X-ray photoelectron spectroscopy · HP-XPS · atomic force microscopy · AFM

Deactivation of catalytic Pt nanoparticles (NPs) plays an important role in heterogeneous catalysis since Pt NPs are active for a number of industrially relevant reactions such as oxidation of CO and hydrocarbons in automotive catalytic converters, NO<sub>x</sub> abatement, volatile organic compound (VOC) combustion, catalytic hydrogenation and reforming, and direct decomposition of alcohols for hydrogen production.<sup>1–4</sup> However, although the smallest Pt NPs may offer desirable properties such as enhanced catalytic activity and selectivity, as well as the economic advantage of minimizing the precious metal loading, short- and long-term stability issues unfortunately limit their applications. The primary routes to deactivation of Pt NPs are through poisoning,

coarsening, dissolution in electrochemical environments, and/or volatilization under gas phase reaction conditions.<sup>5–11</sup> Inherently, NPs are metastable and susceptible to coarsening because the high surface free energy can be minimized by the coarsening process.<sup>12,13</sup> This process is limited kinetically by intraparticle or interparticle diffusion of atomic species.<sup>12</sup> Hence, increased temperatures tend to facilitate sintering. By intraparticle diffusion of atomic species over the surface of the NP, the whole particle can migrate and cause coalescence, that is, fusion with another NP. Alternatively, Ostwald ripening can proceed through interparticle migration of atomic species. Common to both sintering mechanisms is that the catalyst material is preserved on the support

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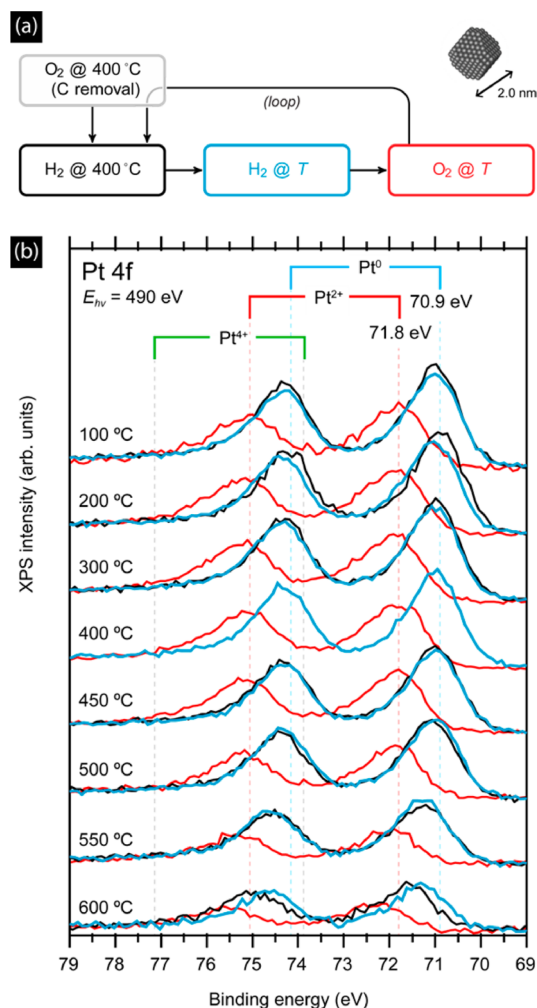
surface, but the catalytic activity is reduced due to the reduced number of active sites (*i.e.*, decreased surface-to-volume ratio).

Recently, the thermal stability of Pt NPs supported on SiO<sub>2</sub> has been the subject of several studies.<sup>13,14</sup> For example, Simonsen *et al.* observed sintering of physical vapor deposited Pt NPs on SiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> under 7.5 Torr of technical air (21% O<sub>2</sub> and 79% N<sub>2</sub>) at 650 °C by *in situ* transmission electron microscopy (TEM), which was attributed to ripening.<sup>13,15</sup> Other previous studies have also discussed accelerated sintering rates in O<sub>2</sub> atmospheres.<sup>12,16</sup> On other supports, such as Al<sub>2</sub>O<sub>3</sub>, enhanced sintering stability and/or dispersion of Pt NPs has been reported after O<sub>2</sub> exposure at lower temperatures ( $\sim$ 400 °C).<sup>17,18</sup> In addition, NP migration and ripening were not observed up to  $\sim$ 700 °C under ultrahigh vacuum (UHV) conditions for micelle-synthesized Pt NPs supported on SiO<sub>2</sub>/Si(111).<sup>14,19</sup>

Here, we report on the stability of Pt NPs synthesized by inverse micelle encapsulation and supported on SiO<sub>2</sub>/Si(111) at conditions closer to those relevant in industrial applications. Using high-pressure XPS, we have studied the properties of the Pt NPs in the presence of elevated pressures of H<sub>2</sub> and O<sub>2</sub> (0.5 Torr). From *in situ* XPS and *ex situ* AFM data, monodispersed 2 nm Pt NPs on a flat SiO<sub>2</sub>/Si(111) substrate show no indication of sintering up to 600 °C in reducing and oxidizing environments. However, it was observed that the stability of catalytically active Pt NPs is hampered by desorption of Pt, presumably as volatile PtO<sub>x</sub> species formed above 450 °C in O<sub>2</sub>. For material systems with limited diffusivity such as NPs dispersed in a porous support, this PtO<sub>x</sub> desorption may result in coarsening of the NPs caused by readsorption of the volatile material.

## RESULTS AND DISCUSSION

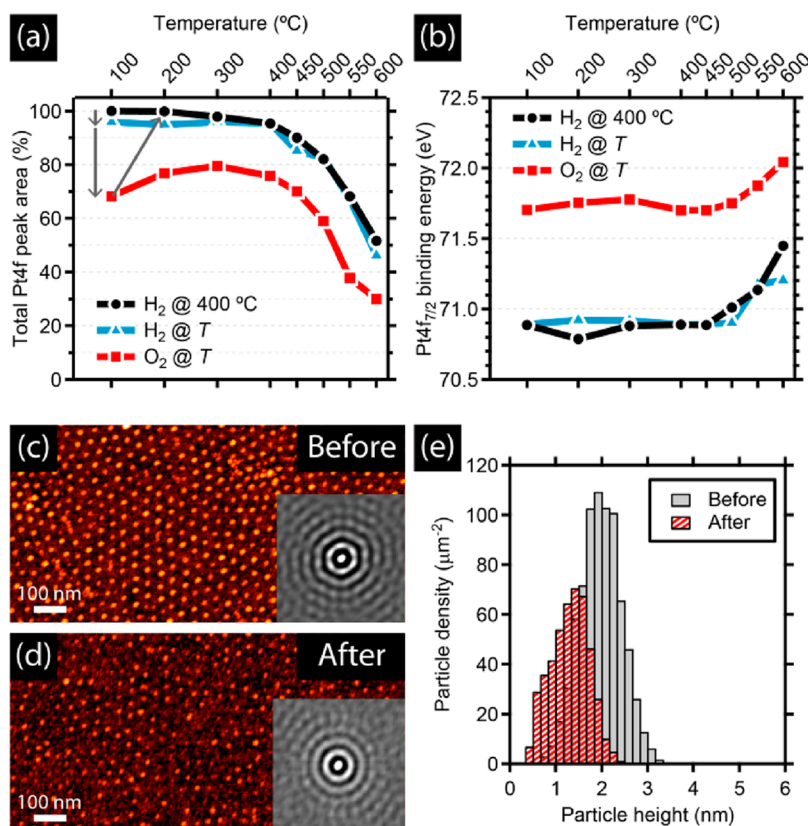
As depicted in Figure 1a,  $\sim$ 2 nm Pt NPs supported on SiO<sub>2</sub>/Si(111), sample S1, were cleaned by annealing in 0.5 Torr O<sub>2</sub> (see also the Experimental Procedure section) followed by reduction in 0.5 Torr H<sub>2</sub> at 400 °C. Subsequently, the temperature was changed to a given temperature (*T*) and measured first in H<sub>2</sub> followed by O<sub>2</sub>. This cycle, beginning with reduction in 0.5 Torr H<sub>2</sub> at 400 °C, was repeated for different temperatures (*T*) from 100 to 600 °C. Figure 1b shows HP-XPS spectra of the Pt 4f core level region. The black and blue curves were recorded in H<sub>2</sub> at 400 °C and at the stated temperatures, respectively, while the spectra depicted in red were recorded in O<sub>2</sub> at the temperatures given. The black spectra show a slight, but clear, increase of the Pt<sup>0</sup> 4f<sub>7/2</sub> BE from 70.9 to 71.4 eV, followed by a decrease in the peak intensities with increasing annealing temperature. The BE values of the Pt 4f<sub>7/2</sub> species above are similar to reported values for metallic Pt.<sup>14,19–23</sup> No noticeable changes are observed in the Pt 4f peak shapes with temperature when the Pt NPs are exposed to reducing conditions,



**Figure 1.** (a) Flow diagram of the HP-XPS measurement steps conducted on sample S1, Pt NPs ( $\sim$ 2 nm) supported on SiO<sub>2</sub>/Si(111). The sample was initially annealed in O<sub>2</sub> for the removal of residual carbon followed by H<sub>2</sub> for NP reduction. Subsequently, the sample was subjected to heating cycles in 0.5 Torr of H<sub>2</sub> and O<sub>2</sub> at specified temperatures (*T*) increasing in steps from 100 to 600 °C. In each cycle, before annealing at a given temperature in either O<sub>2</sub> or H<sub>2</sub>, the sample was measured in 0.5 Torr of H<sub>2</sub> at 400 °C. (b) Series of Pt 4f XPS spectra recorded in 0.5 Torr H<sub>2</sub> at 400 °C (black), 0.5 Torr H<sub>2</sub> at *T* (blue), and 0.5 Torr O<sub>2</sub> at *T* (red). The incident photon energy was 490 eV.

showing a narrow full width half-maximum (fwhm) of 1.1 eV. When the Pt NPs are exposed to O<sub>2</sub>, distinct chemical shifts to higher BEs, ranging from 71.7 to 72.0 eV (Pt 4f<sub>7/2</sub>), are observed. Since the fwhm values determined from fits of our XPS spectra in O<sub>2</sub> ( $\sim$ 1.2 eV) are similar to those in H<sub>2</sub>, the presence of a single chemical state is inferred. The presence of significant amounts of PtO<sub>2</sub> can be ruled out based on the reported Pt<sup>4+</sup> 4f<sub>7/2</sub> BE values from the literature ranging from 73.7 to 74.5 eV.<sup>14,19,24,25</sup> Moreover, if a PtO<sub>2</sub> phase was present, a Pt<sup>4+</sup> 4f<sub>5/2</sub> peak at  $\sim$ 77.1 eV would be expected. Instead, and in agreement with previous reports,<sup>14,19,21,23,26</sup> the oxide species in our samples is assigned to PtO.

Figure 2a,b shows the changes in the integrated areas and BEs of the Pt 4f XPS peaks, respectively, as



**Figure 2.** Temperature dependence of the (a) integrated Pt 4f XPS peak areas and (b) Pt 4f<sub>7/2</sub> XPS binding energies of sample S1, Pt NPs (~2 nm) supported on SiO<sub>2</sub>/Si(111), in 0.5 Torr H<sub>2</sub> at 400 °C (black circles), 0.5 Torr H<sub>2</sub> at T (blue triangles), and 0.5 Torr O<sub>2</sub> at T (red squares). The values for the integrated areas are given in percent relative to the first point measured in 0.5 Torr H<sub>2</sub> at 400 °C. Gray arrows indicate the order of the measurements. *Ex situ* AFM images of Pt NPs (c) before and (d) after the high-pressure experiments. (e) NP height histogram extracted from the AFM images. Insets in (c) and (d) show the corresponding autocorrelation images.

a function of the annealing temperature in H<sub>2</sub> (blue triangles) and O<sub>2</sub> (red squares) at specified temperatures (top horizontal axis). Reference data acquired in H<sub>2</sub> at 400 °C after each treatment are also shown (black circles). During each cycle, when the HP-XPS spectra are measured in the presence of O<sub>2</sub> (red squares), the Pt 4f signal intensities appear smaller as compared to those in H<sub>2</sub> at the same temperature. This is attributed to the attenuation of the signal due to the presence of oxygen atoms in the oxide phase and possibly also to reversible shape changes of the Pt NPs in O<sub>2</sub>.<sup>27</sup> On the contrary, data measured in H<sub>2</sub> at 400 °C (black circles) after the treatment in O<sub>2</sub> indicate that irreversible changes have taken place. A temperature of ~450 °C is found as the threshold for NP changes, where a strong overall decrease in the Pt 4f peak areas followed by an increase in the BEs is observed. The comparison of the two sequential points (black → blue) in each cycle, Figure 2a, where one corresponds to the measurements in H<sub>2</sub> at 400 °C (black circles) and the second to the subsequent increase in annealing temperature (blue triangles) keeping the same H<sub>2</sub> environment (e.g., 400 °C → 450 °C, 400 °C → 500 °C, 400 °C → 550 °C, and 400 °C → 600 °C), shows no significant differences in the total Pt 4f peak areas. This indicates

that the major irreversible changes in the Pt 4f signal are only observed after the Pt NPs are annealed above 450 °C in an O<sub>2</sub> atmosphere.

The origin of the observed intensity attenuation could be assigned to one or more of the following effects: sintering of the Pt NPs, NP shape changes under different gas environments, or loss of Pt from the NPs. In order to elucidate which of these effects contributes the most, *ex situ* AFM measurements of the treated sample S1 ("After") and an identically prepared, untreated sample ("Before") were acquired and analyzed (Figure 2c,d). AFM is suitable not only for measuring the size distribution of the Pt NPs but also for determining their spatial distribution on the support. As a result of the micellar preparation method,<sup>14,28,29</sup> the distribution of Pt NPs exhibits local hexagonal ordering, with uniform interparticle distances determined by the length of the diblock copolymer tails (PS group). The autocorrelation images shown as insets in Figure 2c,d show distinct rings corresponding to the particles' coordination shells, revealing strong local ordering before and after treatment with average NP nearest-neighbor distances of 40 nm in both cases. Thus, sintering of the NPs by particle migration did not occur for our micellar NPs.<sup>4,14,18,19</sup> Figure 2e shows the AFM NP height

histograms of the as-prepared (gray columns) and the treated sample after the HP-XPS measurements (red columns). The NPs in the as-prepared sample had an average height of  $2.0 \pm 0.5$  nm, while those subjected to the high-pressure experiments had an average height of  $1.3 \pm 0.4$  nm. Because NPs smaller than  $\sim 0.5$  nm are difficult to unambiguously identify in the AFM images due to the rough pearl-like structure of the amorphous SiO<sub>2</sub> support, the precise determination of particle count on the sample after the treatments is difficult. A large number of particles in the range of 0–0.5 nm have not been counted in the latter case, and thus the height distributions obtained from the AFM measurements of the sample after treatment are biased toward a higher average value. Considering the uncounted particles, the true mean height is in the range of 0.8–1.3 nm.

These observations reveal that the decreasing XPS areas obtained with increasing annealing temperature above 450 °C can be primarily assigned to the loss of Pt from the NPs, most likely desorbing as PtO<sub>x</sub> during the annealing step in O<sub>2</sub>. The other plausible explanations for the XPS signal attenuation, such as sintering or shape change, are ruled out since both would lead to increased or unchanged NP heights in AFM images. Additionally, the positive Pt 4f BE shifts observed as a function of annealing temperature above 450 °C (Figure 2b) corroborate the finding that the Pt NPs become smaller. This conclusion is based on the well-established size-dependent initial and final state effects, which are known to lead to positive BE shifts.<sup>30</sup>

The attenuation in the HP-XPS signal is consistent with the reduction in NP size measured by AFM. If a NP of height  $h_1$  is reduced to a NP of height  $h_2$ , and with further assumption that the NP has either spherical or hemispherical shape, the ratio of XPS intensities is given (through application of the Beer–Lambert attenuation law) by

$$\frac{I(h_2)}{I(h_1)} = \frac{1 - \left(1 + \frac{h_2}{\lambda}\right) \exp\left(-\frac{h_2}{\lambda}\right)}{1 - \left(1 + \frac{h_1}{\lambda}\right) \exp\left(-\frac{h_1}{\lambda}\right)} \quad (1)$$

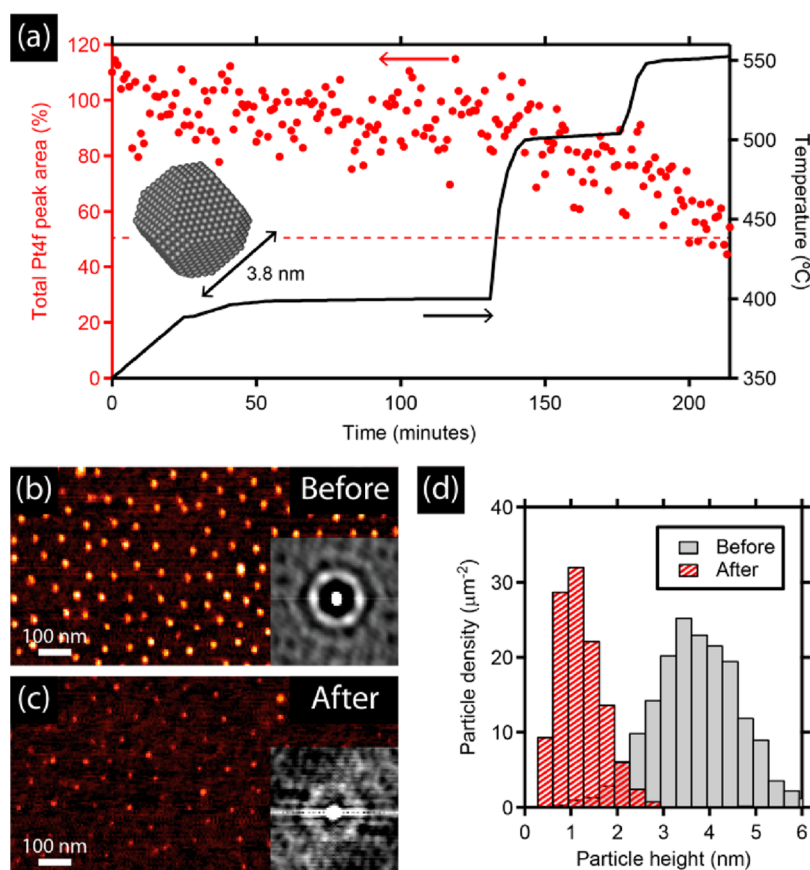
where  $\lambda$  is the inelastic mean free path (IMFP). From the AFM analysis, the initial NP height was  $2.0 \pm 0.5$  nm (Figure 2e).<sup>31</sup> Thus, if we use 2.0 nm as  $h_1$  in eq 1 and  $I(h_2)/I(h_1) = 0.50$ , we can then estimate the final NP size numerically. Using the value of 0.733 nm for the IMFP of electrons with a kinetic energy of 420 eV,<sup>32,33</sup> we find a value of 1.0 nm for  $h_2$ , which is consistent with the measured AFM height of 0.8–1.3 nm.

In order to verify that the observed loss of Pt is mediated by the elevated pressure of O<sub>2</sub>, we have studied the thermal stability of another sample with  $\sim 4$  nm Pt NPs (S2), exclusively in the presence of O<sub>2</sub> after the initial cleaning in O<sub>2</sub> and reduction in H<sub>2</sub> at 400 °C. The time dependence of the XPS Pt 4f peak area (red circles, left axis) and the temperature (black line, right axis) is shown in Figure 3a. A similar onset temperature

of  $\sim 450$  °C for the drastic Pt 4f XPS attenuation was observed, confirming that the attenuation of the Pt signal detected for S1 at high temperatures was indeed caused by the high pressure of O<sub>2</sub>. Figure 3b,c shows AFM images of sample S2 without treatment (“Before”) and after the high-pressure experiments (“After”), respectively. On the basis of the autocorrelation images shown in the insets, the interparticle distance is observed to be practically identical before (69 nm) and after (72 nm) HP-XPS measurements. However, the height histograms shown in Figure 3d show a drastic decrease in the NP height from an initial average height of  $3.8 \pm 1.3$  nm to  $1.3 \pm 0.7$  nm. Using eq 1, a 50% decrease in the HP-XPS signal (Figure 3a) leads to an estimated 1.4 nm final NP size, indicating good agreement between the model and AFM results. The IMFP was set to 0.868 nm, corresponding to 550 eV electrons.<sup>32,33</sup>

The loss of Pt from the NPs can be explained in terms of the Pt oxidation kinetics and thermodynamics. The entire NPs were observed to be oxidized in the presence of O<sub>2</sub>, but with no indication of PtO<sub>2</sub> formation, in agreement with a report by Wang *et al.*<sup>34</sup> for  $\sim 2.0$  nm Pt NPs. It should be however noted that there have also been theoretical reports suggesting PtO<sub>2</sub> as the most thermodynamically stable phase for Pt NPs <6 nm under our experimental conditions.<sup>35,36</sup> This finding suggests that the oxidation to PtO<sub>2</sub> is kinetically hindered by O<sub>2</sub> dissociation on the oxidic Pt NPs or by barriers for intraparticle mass transport.<sup>37</sup> This is in agreement with previous reports that PtO, rather than PtO<sub>2</sub>, was found to be the predominant species after thermal treatments of supported Pt NPs in O<sub>2</sub>,<sup>4,34</sup> and that PtO<sub>2</sub> is observed after atomic oxygen exposure<sup>14,19</sup> where the barrier for dissociative adsorption of O<sub>2</sub> is bypassed. It is plausible that the onset temperatures for PtO<sub>x</sub> sublimation<sup>35,38</sup> are lower than the temperature required to overcome the kinetic barrier for PtO<sub>x</sub> formation. Consequently, the formation of PtO<sub>x</sub> on the NP surface might immediately lead to its desorption. The desorbing species may be PtO<sub>2</sub>.

Pt volatilization has been observed in several studies of Pt catalysts under oxidizing conditions similar to those employed here.<sup>12,39–41</sup> For example, Lamber and Romanowski<sup>40</sup> investigated the stability of Pt particles grown on quartz by physical vapor deposition in vacuum and in low pressures of oxygen ( $2.3 \times 10^{-4}$  Torr) and observed substantial decreases in crystallite size and Pt mass loss following oxygen treatments at 500 °C, while the vacuum treatments led to NP growth. As discussed by Wynblatt and Gjostein<sup>12</sup> and Harris,<sup>42</sup> vaporization of PtO<sub>2</sub> can contribute to NP coarsening in a similar manner to surface diffusion, *via* the exchange of Pt between particles through an Ostwald ripening process. The fact that only a decrease in NP size during O<sub>2</sub> exposure at high temperature was observed in our



**Figure 3.** (a) Temperature and time evolution of the integrated Pt 4f XPS peak areas (closed circles) of sample S2, Pt NPs ( $\sim 4$  nm) supported on  $\text{SiO}_2/\text{Si}(111)$ , in  $\text{O}_2$  atmosphere (0.5 Torr). AFM images of the NPs (b) before and (c) after the high-pressure experiments. (d) Histogram of NP height measurements extracted from the AFM images. Insets in (b) and (c) show the corresponding autocorrelation images. In (a), the XPS spectra were recorded with an incident photon energy of 620 eV.

study and that of Lamber is most likely due to the very low rate of readsorption on our samples due to the planar geometry and the long diffusion length in the gas phase under the conditions of our experiments.

It should be noted that, in the similar studies by Simonsen *et al.* employing *in situ* TEM of Pt NPs on  $\text{SiO}_2$  and  $\text{Al}_2\text{O}_3$ , significant coarsening of the NPs was observed in 7.5 Torr technical air (21%  $\text{O}_2$  and 79%  $\text{N}_2$ ) at 650 °C without apparent loss of Pt from the surfaces.<sup>13,15</sup> Similar results were obtained by Chu and Ruckenstein,<sup>43</sup> who studied Pt evaporated onto an  $\text{Al}_2\text{O}_3$  film, though in their case, NP diffusion and coalescence were found to be the primary mode of sintering. In the HP-XPS experiments reported here, the differential pumping of gas into the analyzer causes substantial gas flow near the sample surface, which might help minimize Pt (or  $\text{PtO}_x$ ) redeposition. In the studies by Simonsen and Chu, it is possible that the higher gas pressures and lower gas flows near the surface allowed such redeposition. Potential differences in NP size, shape, and interaction with the support due to the micellar preparation method (exposure to oxygen plasma) may also play a role, as could the initially substantially larger (40–70 nm vs

$\sim 2$ –5 nm as in refs 13 and 15) interparticle distances in our samples, potentially allowing diffusing surface  $\text{PtO}_x$  species to desorb before finding another particle.

The present work highlights the complexity of the fundamental processes underlying catalyst stability and emphasizes the importance of model catalyst studies in gaining a fundamental understanding of the dynamics of NP oxidation and sintering at the nanoscale. Our study suggests that  $\text{PtO}_x$  volatility plays an important role in catalyst stability under oxidizing conditions.

## CONCLUSION

From an interplay of *in situ* HP-XPS and *ex situ* AFM measurements, the thermal stability of micellar Pt NPs supported on  $\text{SiO}_2/\text{Si}(111)$  was investigated under 0.5 Torr of reducing and oxidizing gaseous ( $\text{H}_2$  and  $\text{O}_2$ ) environments. No indication of NP sintering (Ostwald ripening or diffusion coalescence) was observed up to 600 °C after exposure to  $\text{H}_2$  and  $\text{O}_2$ , although such an effect could be counteracted by the Pt loss observed in  $\text{O}_2$  at temperatures above  $\sim 450$  °C. When probed under oxidizing conditions, HP-XPS shows a drastic decrease in the Pt 4f signal at temperatures above  $\sim 450$  °C. It is shown that this irreversible Pt 4f signal

attenuation is caused by the formation of volatile  $\text{PtO}_x$  species. However, comparison of the AFM autocorrelation images acquired before and after the HP-XPS experiments

did not show any clear changes in the spatial distribution of the particles (*i.e.*, NP mobility), indicating the strong adhesion of the particles to the support.

## EXPERIMENTAL PROCEDURE

Size-selected Pt NPs were synthesized by the inverse micelle encapsulation technique as described in detail in refs 14, 28, and 29. Briefly, two types of commercially available nonpolar/polar diblock copolymers, poly(styrene)-*block*-poly(2-vinylpyridine) [S1: PS(27700)-*b*-P2VP(4300) and S2: PS(53000)-*b*-P2VP(43800)], were dissolved in toluene in order to form inverse micelles. Subsequently, the micelles were loaded with a  $\text{H}_2\text{PtCl}_6 \cdot 6\text{H}_2\text{O}$  salt. The size selection was obtained by varying the molecular weight of the polymer core (P2VP) and by changing the metal salt/P2VP ratio. The metal loadings (metal/P2VP ratio) used were 0.2 for S1 and 0.05 for S2. The interparticle distance was determined by the lengths of the nonpolar tails (PS). A monolayer-thick film of polymer-encapsulated Pt NPs was obtained upon dip-coating naturally oxidized amorphous  $\text{SiO}_2/\text{Si}(111)$  substrates into the metal-loaded polymeric solution. The encapsulating polymer was removed in a commercial UHV surface analysis system (SPECS GmbH) via an  $\text{O}_2$  plasma treatment ( $P[\text{O}_2] = 3.0 \times 10^{-5}$  Torr for 100 min) at room temperature. Samples with NP sizes of  $2.0 \pm 0.5$  nm (S1) and  $3.8 \pm 1.3$  nm (S2) were obtained by the method described above.

HP-XPS experiments were performed at the beamline 9.3.2 of the Advanced Light Source (ALS) at Lawrence Berkeley National Laboratory.<sup>44,45</sup> The system consists of an analysis chamber with an electron spectrometer mounted on a differentially pumped electrostatic lens system that allows XPS measurements in gas environments up to 1 Torr. The base pressure in this chamber is  $\sim 3 \times 10^{-10}$  Torr.  $\text{O}_2$  and  $\text{H}_2$  gases were dosed individually by backfilling the chamber through a leak valve. Each sample transferred to the HP-XPS analysis chamber was further cleaned *in situ* by annealing in an  $\text{O}_2$  atmosphere (0.5 Torr) at 400 °C for 30 min to remove any residual carbon species. After this cleaning procedure, no C 1s signal was observed in XPS. Subsequently, the oxidized Pt NPs were annealed in  $\text{H}_2$  (0.5 Torr) at 400 °C, as shown in Figure 1a. After this treatment, the reduction of the NPs was confirmed by XPS. No significant decrease in the Pt 4f intensity was observed after the cleaning and reduction procedures. Samples S1 and S2 received different treatments during the HP-XPS measurements. As depicted in Figure 1a, sample S1 was alternately measured in  $\text{H}_2$  (0.5 Torr) and in  $\text{O}_2$  (0.5 Torr) for 15 min while keeping the sample at a desired temperature. Prior to each cycle, the sample was annealed at 400 °C in 0.5 Torr of  $\text{H}_2$  until XPS showed complete reduction of the  $\text{PtO}_x$  species to the metallic state. The cycle was repeated by stepwise increasing the temperature for each cycle from 100 to 600 °C, which was the maximum temperature allowed in the experimental setup. Sample S2 was measured only in an  $\text{O}_2$  atmosphere (0.5 Torr), monitoring both the temperature and temporal dependence on the Pt 4f signal decay.

The binding energy (BE) scale was calibrated using the  $\text{Si}^{2p}$  peak at 99.3 eV of the NP substrate [ $\text{SiO}_2(4 \text{ nm})/\text{Si}(111)$ ]. The BE of  $\text{Si}^{2p}$  was measured initially using as reference the Au  $4f_{7/2}$  (84.0 eV) signal of a small Au foil mounted on the sample holder adjacent to the sample. Quantitative analysis was performed by subtracting a Shirley-type background<sup>46</sup> and fitting each XPS spectrum using two asymmetric Voigt line shapes separated by 3.3 eV corresponding to the spin-orbit splitting of Pt  $4f_{7/2}$  and Pt  $4f_{5/2}$  peaks.

Characterization of the sample morphology (NP size and interparticle distance) before and after HP-XPS was performed by *ex situ* AFM imaging in tapping mode (Digital Instruments, Multimode). Noncontact silicon AFM tips with an end radius below 10 nm were used. The cantilevers used have resonant frequencies in the range of 200–300 kHz with spring constants between 20 and 30 N/m. A typical scan rate of 1 Hz was used to acquire the AFM images, and the tip-sample interaction was minimized to avoid tip-induced distortions of the NP morphology

and NP lateral displacements during scanning. Due to tip convolution effects, the NP height (rather than the diameter) has been used here as representative size parameter.

Beam-induced effects were excluded by changing the sample position at the end of each XPS experiment to confirm that the samples remained homogeneous. The beam spot was  $0.5 \text{ mm} \times 1 \text{ mm}$ , while the sample size was  $9 \text{ mm} \times 9 \text{ mm}$ . Thus, beam-induced effects would be easily detected by this procedure. AFM imaging of different areas of the sample showed no indication of inhomogeneity before and after the HP-XPS experiments.

**Conflict of Interest:** The authors declare no competing financial interest.

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